



Mustansiriayah University / College of Science Physics Department Subject: Advance Optics Lecture (4) for Ph.D.

EMW Propagation in dielectrics and conductors

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Plane waves in lossy media: $(\sigma \neq 0)$

$$\vec{\nabla} \times \vec{E} = -j\omega\mu \vec{H}$$

$$\vec{\nabla} \cdot \vec{E} = \frac{\rho}{\epsilon}$$

$$\vec{\nabla} \cdot \vec{B} = 0$$

$$\vec{\nabla} \times \vec{H} = (\sigma + j\omega \epsilon)\vec{E} = j\omega\epsilon_{c}\vec{E}$$

$$\epsilon_{c} = \epsilon - j\frac{\sigma}{\omega} = \epsilon' - j\epsilon'' \qquad including \ damping \ \&ohmic \ losses$$

$$Loss \ tangent \ \tan\theta = \frac{\sigma}{\epsilon w} = \frac{\epsilon''}{\epsilon'}$$

$$\sigma \gg \epsilon \omega \qquad good \ conductor$$

 $\sigma \ll \epsilon \omega$ good insulator





Wave propagation in low-loss dielectrics:

 $tan\theta = \frac{\epsilon^{\prime\prime}}{\epsilon^{\prime}} = \frac{\sigma}{\epsilon\omega} \ll 10^{-2}$ dielectric

For a low-loss dielectric (like ordinary imperfect insulators)

$$\begin{split} \frac{\sigma}{\epsilon\omega} &\ll 1 \quad , \ \epsilon'' \ll \ \epsilon' \qquad \varepsilon_c \ = \epsilon' - j\epsilon'' = \ \epsilon - j\frac{\sigma}{\omega} \\ \gamma &= \alpha + j\beta = j\omega\sqrt{\mu\epsilon'} \ (1 - j\frac{\epsilon''}{\epsilon'})^{1/2} \\ \gamma &= j\omega\sqrt{\mu\epsilon'}(1 - j\frac{\epsilon''}{2\epsilon'} + \frac{1}{8}(\frac{\epsilon''}{\epsilon'})^2) \\ \alpha &= \frac{\omega\epsilon''}{2}\sqrt{\frac{\mu}{\epsilon'}} = \frac{\sigma}{2}\sqrt{\frac{\mu}{\epsilon}} \qquad NP/m \\ \beta &= \omega\sqrt{(\mu\epsilon'}\left(1 + \frac{1}{8}\left(\frac{\epsilon''}{\epsilon'}\right)^2\right) = \ \omega\sqrt{(\mu\epsilon}\left(1 + \frac{1}{8}\left(\frac{\sigma}{\epsilon\omega}\right)^2\right) \\ \eta_c &= \sqrt{\frac{\mu}{\epsilon'}}\left(1 + j\frac{\epsilon''}{2\epsilon'}\right) = \sqrt{\frac{\mu}{\epsilon}}\left(1 + j\frac{\sigma}{2\epsilon\omega}\right) \end{split}$$

Phase velocity

$$v_{\rm p} = \frac{\omega}{\beta} = \frac{1}{\sqrt{(\mu\epsilon')} \left(1 + \frac{1}{8} \left(\frac{\epsilon''}{\epsilon'}\right)^2\right)}$$

Wave propagation in good conductors:

 $\tan \theta = \frac{\epsilon''}{\epsilon'} = \frac{\sigma}{\epsilon_0} \gg 10^2$ Perfect conductor **Good conductor** $\frac{\sigma}{\epsilon \omega} \gg 1 \& \epsilon'' \gg \epsilon'$ $\gamma = \alpha + j\beta = \omega \sqrt{\mu\epsilon} (1 - j\frac{\sigma}{c\omega})^{1/2}$ $\gamma = \sqrt{j}\sqrt{\omega\mu\sigma} = \frac{1+j}{\sqrt{2}}\sqrt{\omega\mu\sigma}$ $\gamma = (1 + j)\sqrt{\pi f\mu\sigma}$ When $\sqrt{j} = (e^{\frac{i\pi}{2}})^{\frac{1}{2}} = e^{\frac{i\pi}{4}} = \cos(\frac{\pi}{4}) + j\sin(\frac{\pi}{4}) = \frac{1+j}{\sqrt{2}}$ $\alpha = \beta \cong \sqrt{\pi f \mu \sigma} \qquad good \ conductor$ $\eta_{\rm c} = \sqrt{\frac{\mu}{\epsilon'}} \left(1 - j\frac{\epsilon''}{\epsilon'}\right)^{-1/2} = \sqrt{\frac{j\omega\mu}{\sigma}} = (1+j)\sqrt{\frac{\pi f\mu}{\sigma}}$ $\eta_c = (1+j)\frac{\alpha}{r}$ $\vec{E} = \hat{i}E_0e^{-\alpha z}e^{-j\beta z}$ $\vec{H} = \frac{E}{\eta_{c}} = \hat{j} \frac{E_{0}}{\eta_{c}} e^{-\alpha z} e^{-j\beta z} = \hat{j} \frac{E_{0}}{\sqrt{\frac{j\omega\mu}{\sigma}}} e^{-\alpha z} e^{-j\beta z}$ $\vec{H} = \hat{j} \frac{E_0}{\sqrt{\frac{\omega\mu}{\sigma}}} e^{-\alpha z} e^{-j(\beta + \frac{\pi}{4})} \qquad H(z) \text{ lags behind } E(z) \text{ by } \frac{\pi}{4}$

Phase velocity

$$v_{\rm p} = \frac{\omega}{\beta} = \sqrt{\frac{2\omega}{\mu\sigma}} \ll c$$
$$v_{\rm p} \downarrow \equiv \sigma \uparrow$$

Example 1: For copper (cu) with $\sigma = 5.8 \times 10^7 S/m$, and phase velocity ($v_p = 720 m/s$) at f = 3 MHz: $\mu = \mu_0 = 4\pi * 10^{-7}$ H/m.

Wavelength:

$$\lambda = \frac{2\pi}{\beta} = \frac{\nu_p}{f} = 2\sqrt{\frac{\pi}{f\mu\sigma}} = 0.24$$
mm

$$\lambda \downarrow = \sigma \uparrow$$

For $\lambda = 0.24$ mm $\ll 100$ m in air at f = 3 MHz. Where $\lambda = \frac{c}{f} = 100$ m

Skin depth δ = Depth of penetration of a good conductor

= Distance thru which the wave amplitude decrease by (e^{-1})

For
$$\vec{E} = \hat{i}E_0 e^{-\gamma z} = \hat{i}E_0 e^{-\alpha z} e^{-j\beta z}$$



$$\delta = \frac{1}{\alpha} = \frac{1}{\sqrt{\pi f \mu \sigma}}$$
$$\delta = \frac{1}{\beta} = \frac{\lambda}{2\pi}$$

 $\delta\downarrow$ as $\sigma\uparrow$ and/or f \uparrow

Cu
$$\delta = 0.038 \text{ mm}$$
 at $f = 3 \text{ MHz}$
 $\delta = 0.66 \mu \text{m}$ at $f = 10 \text{ GHz}$
 $\delta = \frac{1}{\alpha} = \frac{1}{\sqrt{\pi f \mu \sigma}} = \frac{1}{(\pi * 50 * 4\pi * 10^{-7} * 5.8 * 10^{-7})^{\frac{1}{2}}} = ?? \text{m}$ at $f = 50 \text{ Hz}$

Example 2: A LP plane wave $\vec{E} = \hat{\iota}E(z,t)$ propagating along +z-direction in seawater($\epsilon_r = 72, \mu_r = 1, \sigma = 4 S/m$) with $E_0 = \hat{\iota}100 \cos(10^7 \pi t)$ V/m at z=0. Find

a) $\alpha, \beta, \eta_c, v_p, \lambda \& \delta$ b) Z_1 when $E_1 = 0.01E_0$ c) E(0.8, t) & H(0.8, t)

Solutions:

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$$\omega = 10^{7} \pi$$

$$f = \frac{\omega}{2\pi} = 5 \times 10^{6} Hz$$

$$tan\theta = \frac{\sigma}{\epsilon\omega} = 200 \gg 1$$
good conductor
a) $\alpha = \beta = \sqrt{\pi f \mu \sigma} = 8.89 \ rad/m$
 $\eta_{c} = (1+j) \sqrt{\frac{\pi f \mu}{\sigma}} = \pi e^{\frac{j\pi}{4}}$
 $v_{p} = \frac{\omega}{\beta} = 3.53 \times 10^{6} \ m/s$
 $\lambda = \frac{2\pi}{\beta} = 0.707 \ m$
 $\delta = \frac{1}{\alpha} = 0.112 \ m$
b) $E_{1} = 0.01E_{0}$
 $E_{1} = E_{0}e^{-\alpha z_{1}}$
 $0.01E_{0} = E_{0}e^{-\alpha z_{1}}$
 $-\alpha z_{1} = \ln (0.01)$

$$z_1 = \frac{-\ln (0.01)}{\alpha} = 0.518 \, m$$

c)
$$E(z) = \hat{1}100e^{-\alpha z} e^{-j\beta z}$$
 in the phasor domain
 $E(z,t) = R_e(E(z)e^{j\omega t}) = \hat{1}100e^{-\alpha z}\cos(\omega t - \beta z)$

 $\therefore E(0.8, t) = \hat{i} \ 0.082 \cos(10^7 \ \pi t - 7.11)$

$$H(z,t) = R_{e} \left(\hat{j} \frac{E(z)}{\eta_{c}} e^{j\omega t} \right) = R_{e} \left(\hat{j} \frac{100e^{-\alpha z} e^{-j\beta z}}{\pi e^{\frac{j\pi}{4}}} e^{j\omega t} \right)$$

$$H(z,t) = R_{e} \left(\hat{j} \frac{100}{\pi} e^{-\alpha z} e^{j(\omega t - \beta z - \frac{\pi}{4})} \right)$$

$$H(z,t) = \hat{j} \left(\frac{100}{\pi} e^{-\alpha z} \cos(\omega t - \beta z - \frac{\pi}{4}) \right)$$

$$H(0.8,t) = \hat{j} \left(\frac{100}{\pi} e^{-0.8\alpha} \cos(\omega t - 0.8\beta - \frac{\pi}{4}) \right) \text{ ignoring } \omega t \text{ term}$$

$$H(0.8,t) = \hat{j} \ 0.026 \cos(10^{7} \pi t - 7.89)$$

$$H(0.8,t) = \hat{j} \ 0.026 \cos(10^{7} \pi t - 2\pi - 1.61)$$

$$H(0.8,t) = \hat{j} \ 0.026 \cos(10^{7} \pi t - 1.61) \quad A/m$$

Note:
$$H = \frac{E}{\eta_c}, \quad \eta_c = \frac{E}{H}$$

EMW- Propagation states:

Free space $\sigma = 0, \ \epsilon = \epsilon_0, \ \mu = \mu_0$ $\alpha = \omega \sqrt{\frac{\mu \epsilon}{2}} (\sqrt{(1 + (\frac{\sigma}{\omega \epsilon})^2 - 1)})$ $\alpha = 0$ $\beta = \omega \sqrt{\epsilon_0 \mu_0} = \frac{\omega}{c}$ $\eta = \sqrt{\frac{j \omega \mu}{\sigma + j \omega \epsilon}}$ $\eta = \sqrt{\frac{\mu_0}{\epsilon_0}}$ $v_p = \frac{\omega}{\beta} = \frac{1}{\sqrt{\epsilon_0 \mu_0}} = c$

E & H in phase & amplitude does not decay

♦ Lossless dielectric medium $\sigma = 0$, $\epsilon \& \mu$

$$\alpha = 0$$

$$\beta = \omega \sqrt{\epsilon \mu} = \frac{\omega}{\nu}$$

$$\eta = \sqrt{\frac{\mu}{\epsilon}} = \sqrt{\frac{\mu_0 \mu_r}{\epsilon_0 \epsilon_r}} = \eta_0 \sqrt{\frac{\mu_r}{\epsilon_r}}$$

$$\nu = \frac{\omega}{\beta} = \frac{\omega}{\omega \sqrt{\mu \epsilon}} = \frac{1}{\sqrt{\mu \epsilon}} = \frac{1}{\sqrt{\mu_0 \mu_r \epsilon_0 \epsilon_r}}$$

$$\nu = \frac{c}{\sqrt{\mu_r \epsilon_r}}$$

E & H in phase & amplitude still does not decay

***** Lossy media: $(\sigma \neq \mathbf{0})$

$$\alpha = \omega \sqrt{\frac{\mu\epsilon}{2}} \left(\sqrt{\left(1 + \left(\frac{\sigma}{\omega\epsilon}\right)^2 - 1\right)} \right)$$
$$\beta = \omega \sqrt{\frac{\mu\epsilon}{2}} \left(\sqrt{\left(1 + \left(\frac{\sigma}{\omega\epsilon}\right)^2 - 1\right)} \right)$$
$$\eta = \sqrt{\frac{j\omega\mu}{\sigma + j\omega\epsilon}}$$
$$\nu = \frac{\omega}{\beta} = \frac{1}{\sqrt{\frac{\mu\epsilon}{2}} \left(\sqrt{\left(1 + \left(\frac{\sigma}{\omega\epsilon}\right)^2 - 1\right)} \right)}$$
$$\nu = \sqrt{\frac{2}{\mu\epsilon} \left(\sqrt{\left(1 + \left(\frac{\sigma}{\omega\epsilon}\right)^2 - 1\right)} \right)}$$

E & H out of phase & amplitude decay

$$\beta = \sqrt{\frac{\omega\mu\sigma}{2}} = \sqrt{\pi f\mu\sigma}$$

Skin depth $(\delta = \frac{1}{\alpha})$

 $\eta = \sqrt{\frac{j\omega\mu}{\sigma}}$ $|\eta| = \sqrt{\frac{\omega\mu}{\sigma}}$ $\arg(\eta) = \frac{\pi}{4}$

E & H are out of phase

Very strong attenuation, wave tend to reflect from good conductor so often do not experience the loss & E leads H by $\frac{\pi}{4}$.

Refractive index (dielectric):

The refractive index of dielectric medium given by:

 $n \geq 1 \quad \& \quad v < c$

For non-magnetic material ($\mu_r = 1$)

Power flow & energy density:

$$\frac{U_{e}}{U_{m}} = \frac{\frac{1}{2} \in E^{2}}{\frac{1}{2} \mu H^{2}} = \frac{\epsilon E^{2}}{\mu H^{2}} = \frac{\epsilon}{\mu} \eta^{2} \qquad \text{where } \eta = \sqrt{\frac{\mu}{\epsilon}}$$

$$\langle s \rangle = \frac{E_{rms}^2}{\eta}$$
 Energy flow

Refraction of E & µ fields:

To determine the refraction of E-field across the interface between two mediums, used boundary conditions (B. C. S)



<u>B.C.S.</u>

$$E_{t1} = E_{t2} \quad \dots \quad (1)$$

$$D_{n1} = D_{n2} \quad \dots \quad (2)$$

$$E_1 \sin \theta_1 = E_2 \sin \theta_2 \quad \dots \quad (3)$$

$$D_1 \cos \theta_1 = D_2 \cos \theta_2 \quad \dots \quad (4)$$

$$\in_1 E_1 \cos \theta_1 = \in_2 E_2 \cos \theta_2 \quad \dots \quad (5)$$
From eq (5) & eq (3) get

 $\frac{\tan\theta_1}{\tan\theta_2} = \frac{\epsilon_{r1}}{\epsilon_{r2}}$

Law of refraction of the E-field at the boundary

free of charge ($\rho = 0$)

Now used B.C.S to determine the refraction of H- field a cross the interface between two mediums.



<u>B.C.S</u>

 $H_{t1} = H_{t2} \dots \dots (1)$ $B_{U1} = B_{U2} \dots \dots (2)$ $B_1 \cos \theta_1 = B_2 \cos \theta_2 \dots \dots (3)$ $H_1 \sin \theta_1 = H_2 \sin \theta_2 \dots \dots (4)$ $\frac{B_1}{\mu_1} \sin \theta_1 = \frac{B_2}{\mu_2} \sin \theta_2$ From eq. (5) & eq. (3) get

 $\frac{\tan\theta_1}{\tan\theta_2} = \frac{\mu_1}{\mu_2} = \frac{\mu_{r1}}{\mu_{r2}}$

Law of refraction of the magnetic flux lines of the boundary with no current density flow (J = 0)

<u>The relation between $\vec{E} \& \vec{H}$:</u>

 $\widehat{u_E}$: unit vector along \vec{E}

 $\widehat{u_H}$: unit vector along \overrightarrow{H}

 $\overrightarrow{\vec{E} \times H} = \vec{S}$ Pointing vector

Pointing vector is the same direction of propagation $(\widehat{u_S})$ or $(\widehat{u_K})$:

$$\widehat{u_{E}} \times \widehat{u_{H}} = \widehat{u_{S}} = \widehat{u_{K}} \dots \dots \dots (1)$$

$$\widehat{u_{K}} \times \widehat{u_{H}} = -\widehat{u_{E}} \dots \dots \dots (2)$$

$$\widehat{u_{K}} \times \widehat{u_{E}} = \widehat{u_{H}} \dots \dots \dots (3)$$

$$\overrightarrow{H} = \frac{1}{\eta} \ \widehat{u_{K}} \times \overrightarrow{E} \dots \dots \dots (4)$$

$$\overrightarrow{E} = -\eta \ \widehat{u_{K}} \times \overrightarrow{H} \dots \dots \dots (5)$$

If a uniform plan wave travel in the +z-direction may have x- & ycompounds:

$$\vec{\mathrm{E}}_{(z)} = \hat{\imath}\vec{\mathrm{E}}_{x}^{+}(z) + \hat{\jmath}\vec{\mathrm{E}}_{y}^{+}(z) \, \dots \, \dots \, (6)$$

The associated H- field is

$$\vec{\mathrm{H}}_{(z)} = \hat{\imath} \vec{\mathrm{H}}_{x}^{+}(z) + \hat{\jmath} \vec{\mathrm{H}}_{y}^{+}(z) \, \dots \, \dots \, (7)$$

The exact expression of magnetic field in terms of electrical field will be:

$$\vec{\mathrm{H}}_{(z)} = \frac{1}{\eta}\hat{k} \times \vec{\mathrm{E}}_{(z)} = -\hat{\imath}\frac{\vec{\mathrm{E}}_{\mathcal{Y}}^{+}(z)}{\eta} + \hat{\jmath}\frac{\vec{\mathrm{E}}_{\mathcal{X}}^{+}(z)}{\eta}$$

$$\vec{\mathrm{H}}_{\mathrm{x}}^{+}(\mathrm{z}) = -\frac{\vec{\mathrm{E}}_{y}^{+}(\mathrm{z})}{\eta}$$
$$\vec{\mathrm{H}}_{\mathrm{x}}^{+}(\mathrm{z}) = \frac{\vec{\mathrm{E}}_{x}^{+}(\mathrm{z})}{\eta}$$

Where
$$\widehat{\mathbf{u}_{\mathrm{S}}} = \widehat{\mathbf{u}_{\mathrm{K}}} = \widehat{k}$$

Summary:

	Any medium	Lossless medium $\sigma = 0$	Low-loss medium $rac{\epsilon''}{\epsilon'}\ll 1$	$\begin{array}{c} Good\\ conductor\\ \frac{\epsilon^{\prime\prime}}{\epsilon^{\prime}} \gg 1 \end{array}$	Unit
α =	$\omega \sqrt{\frac{\mu \epsilon^{'}}{2}} (\sqrt{(1 + (\frac{\epsilon^{\prime\prime}}{\epsilon^{'}})^2} - 1)$	0	$\frac{\sigma}{2}\sqrt{\frac{\mu}{\epsilon}}$	$\sqrt{\pi f \mu \sigma}$	$\frac{Np}{m}$
β =	$\omega \sqrt{\frac{\mu \epsilon'}{2} (\sqrt{(1 + (\frac{\epsilon''}{\epsilon'})^2 - 1)})}$	ω $\sqrt{\mu\epsilon}$	$ω\sqrt{\mu\epsilon}$	$\sqrt{\pi f \mu \sigma}$	$\frac{rad}{m}$
η _c =	$\sqrt{\frac{\mu}{\epsilon'}} \left(1 - j\frac{\epsilon''}{\epsilon'}\right)^{-1/2}$	$\sqrt{\frac{\mu}{\epsilon}}$	$\sqrt{\frac{\mu}{\epsilon}}$	$(1+j)\frac{\alpha}{\sigma}$	Ω
<i>v</i> _p =	$\frac{\omega}{\beta}$	$\frac{1}{\sqrt{\mu\epsilon}}$	$\frac{1}{\sqrt{\mu\epsilon}}$	$\sqrt{\frac{2\pi f}{\mu\sigma}}$	$\frac{m}{s}$
λ =	$\frac{2\pi}{\beta} = \frac{v_{\rm p}}{f}$	$\frac{v_{\rm p}}{f}$	$\frac{v_{\rm p}}{f}$	$\frac{v_{\rm p}}{f}$	m

Notes:

$$\epsilon' = \epsilon \quad \& \quad \epsilon'' = \frac{\sigma}{\omega}$$
$$\epsilon' = \frac{\beta^2 - \alpha^2}{\omega^2 \mu}$$
$$\epsilon'' = \frac{2\alpha\beta}{\omega^2 \mu}$$

In free space $\epsilon = \epsilon_0$, $\mu = \mu_0$ Low-loss medium $\frac{\epsilon''}{\epsilon'} = \frac{\sigma}{\omega\epsilon} < 0.01$ Good conductor $\frac{\epsilon''}{\epsilon'} = \frac{\sigma}{\omega\epsilon} > 100$ Loss tangent $\tan\theta = \frac{\epsilon''}{\epsilon'} = \frac{\sigma}{\epsilon\omega}$ In general $\eta_c = \sqrt{\frac{\mu}{\epsilon'}} (1 - j\tan\theta)^{-1/2}$

dB scale

Power intensity ratio in log scales not a unit:

Attenuation:

$$E_{(Z)} = E_0 e^{-\gamma z} = E_{0 e^{-\alpha z}} e^{-j\beta z}$$

$$A_{(Z)} = 20 \log \left(\frac{E_{(2)}}{E_{(0)}}\right) = 20 \log|e^{-\alpha z}|$$

$$A_{(z)} = \frac{-20 \alpha z}{\ln(10)} dB$$

Example3: If E-field intensity going through a medium attenuates at a rate of (0.4 dB/m) what is a.

Solution:

$$-0.4 = -8.686 \ \alpha + 1m$$
$$\alpha = \frac{0.4}{8.686} = 0.046 \frac{nepers}{m}$$

Note: α positive number for attention.

Example4: The sinusoidal electric field with $E_0 = 250 \text{ v/m}$ and frequency f=1 G Hz exists in a lossy dielectric medium with $E_r = 2.5$ and loss tangent of 0.001. Find the average power dissipated in the medium per cubic meter.

Solution:

$$\tan \theta = 0.001 = \frac{\sigma}{\omega \epsilon} = \frac{\sigma}{\omega \epsilon_0 \epsilon_r} = \frac{\sigma}{2 \pi x \, 2 \, x \, 10^9 \frac{10^{-9}}{36\pi} x \, 2.5}$$

 $\sigma = 1.39 \ge 10^{-4} \ s/m$

The average power P_{ave} dissipated per unit volume V:

$$\frac{P_{ave}}{V} = \frac{1}{2}\vec{J}.\vec{E} = \frac{1}{2}\sigma E^2 = \frac{1}{2}(1.39 \ x \ 10^{-4}) \ x \ (250)^2$$

 $= 4.34 \text{ w/m}^2$

Note:

$$P_{ave} = \frac{1}{2} \frac{v^2}{R} = \frac{1}{2} \frac{(El)^2}{2\rho l/A} = \frac{1}{2} \sigma E^2(lA)$$

v:volt, R:electrical resistance of the uniform specimen

l : length of the specimen A : Area cros – section of the specimen ρ: electrical resistivity

Skin depth:

 $\delta = \frac{1}{\alpha} \sqrt{\frac{2}{\mu \sigma \omega}} \quad \dots \quad (1$ E (z) = E₀ $e^{-\alpha z} e^{-j\beta z} \quad \dots \quad (2$ At z= $\delta \implies |E|$ decreases to $(\frac{1}{e})$ or (0.63 drop) A (z) = 20 log $|\frac{E(z)}{E(0)}| = 20 \log |e^{-\alpha z}|$ A (z) = - 8.686 $\propto z$ dB at z = δ , |E| decrease by (-8.7) dB at z = 2 δ , |E| decrease by (-17.3) dB.

Example5: The skin depth or non-magnetic conducting medium is $(2\mu m)$ at f=5 GHz, Find phase velocity v_p in this medium, then find the attenuation in dB, when the wave penetrates $(10 \ \mu m)$ into the material? Solution:

 $v_{p} = \frac{\omega}{\beta}$ For conductor $\propto = \beta = \frac{1}{\delta}$ $v_{p} = \omega \delta = 2\pi \text{ x } 5 \text{ x } 10^{9} \text{ x } 2 \text{ x } 10^{-6} = 6.28 \text{ x } 10^{4} \text{ m/s}$ A (z) = 20 log $|\frac{E(z)}{E(0)}| = 20 \log |e^{-\alpha z}| = -8.686 \propto z$ A (z) = -8.686 z/ δ = -8.686 . $\frac{10}{2}$ = -43.4 dB high loss There are only surface current on conductors.

Example6:

- a) Calculate the diclectric loss (in dB) of an_EM ware propagating through (100 m) of teflon at f=(MHz)
- b) at f = 10 GHz. $\epsilon_r r = 2.08$, tan $\theta = 0.0004$ at 25 c^0 assuming frequency independeue

Solution:

a) $\tan\theta = \frac{\sigma}{\omega\epsilon}$ $\sigma = \omega\epsilon_0 \epsilon_r \tan\theta = 2\pi \times 10^6 \times \frac{10^{-9}}{36\pi} 2.08 \times 0.0004$ $\sigma = 4.6 \times 10^{-8} \text{ s/m}$ $\alpha = \frac{\sigma}{2} \sqrt{\frac{\mu}{\epsilon}} = \frac{\sigma \eta_{\circ}}{2\sqrt{\epsilon_r}} = \frac{4.6 \times 10^{-8} \times 377}{2 \times \sqrt{2.08}}$ $\alpha = 6.04 \times 10^{-6} \text{ Np/m}$ A (z) = -8.686 \approx z = -8.686 \times 6.04 \times 6.04 \times 10^{-6} \times 100 = - 0.005 dB

b)
$$\sigma = \omega \epsilon_0 \epsilon r \tan \theta = 4.6 \times 10^{-4} \text{ s/m}$$

 $\alpha = \frac{\sigma}{2} \sqrt{\frac{\mu}{\epsilon}} = \frac{\sigma \eta_\circ}{2\sqrt{\epsilon_r}} = 6.04 \times 10^{-2} \text{ Np /m}$
A (z) = - 8.68 \approx z = -8.68 \pm 6.04 \pm 10^{-2} \pm 100
= - 50 dB

Coaxial cable works well at low frequency (Tv to antenna) but not so well at high frequency.

Example 7: In a non – magnetic , lossy , dielectric medium a plane was of frequency f=30 MHZ characterized by the magnetic field phasor :

 $H=(\hat{\imath}-j4\hat{k}) e^{-2y} e^{-j9y}A/m$. Determine the time domain expressions for the electric and magntic field vectors.

Solution:

$$\begin{split} H(\mathbf{r},\mathbf{t}) &= \mathbf{R}_{e} \ (\hat{\imath} - \mathbf{j}4\hat{k}) \ e^{-2y} \ . \ e^{\mathbf{j}(wt-9y)} \ . \\ H(\mathbf{r},\mathbf{t}) &= \hat{\imath}e^{-2y} \cos (\omega t - 9y) + \hat{k}4e^{-2y} \sin (wt - 9y) \\ &\propto = 2 \ , \ \& \ \beta = 9 \\ &- \omega^{2}\mu \ \overline{e} = \alpha^{2} - \beta^{2} \\ &\omega^{2}\mu \ \overline{e} = 2 \propto \beta \\ &\overline{e} = \frac{2\alpha\beta}{\beta^{2}-\alpha^{2}} = \frac{2X \ 2 \ X^{9}}{9^{2}-2^{2}} = 0.468 = \tan\theta \\ &\epsilon_{r} = \frac{\overline{e}}{\epsilon_{0}} = \frac{\beta^{2}-\alpha^{2}}{\alpha^{2}\mu_{0}\epsilon_{0}} = \frac{77C^{2}}{\alpha^{2}} = \frac{77 \ X \ (3 \ X \ 10^{8}))^{2}}{(2 \ X \ 300 \ X \ 10^{6})^{2})} \\ &\epsilon_{r} = 1.95 \\ &\eta_{c} = \sqrt{\frac{\mu}{\epsilon}} \ (1-\mathbf{j} \ \tan\theta)^{-1/2} = \frac{\eta_{o}}{\sqrt{\epsilon_{r}}} \ (1-\mathbf{j} \ \tan\theta)^{-1/2} \\ &\eta_{c} = \frac{377}{\sqrt{1.95}} \ (1-\mathbf{j} \ 0.468)^{-1/2} = 257 \ e^{j \ 0.22} \\ &\mathbf{E} = \hat{k} \ \eta_{c} \cos (\omega t - 9y) - \hat{\imath} \ 4 \ \eta_{c} \ e^{-2y} \sin (\omega t - 9y) \\ &= \hat{k} 257 \ e^{-2y} \cos (\omega t - 9y + 0.22) - \hat{\imath} \ 1028 \ e^{-2y} \sin (\omega t - 9y + 0.22). \end{split}$$

Example8:

A uniform plane wave propagate in a lossless dielectric in the ⁺z direction. The electric field is given by: $\overrightarrow{E(z,t)} = 377 Cos(\omega t - \frac{4\pi}{3}z + \frac{\pi}{6})\hat{i}$ V/m

The average power density measured was 377 W/m^2 , Find:

(i) Dielectric constant of the material if $\mu = \mu_0$

(ii) Wave frequency

(iii) Magnetic field equation

Solution:

(i) Average power:

$$P_{ave} = \frac{1}{2} \frac{E^2}{\eta} = 377$$
$$\frac{1}{2} \frac{(377)^2}{\eta} = 377$$
$$\to \eta = 377 / 2 = 188.5\Omega$$

For lossless dielectric:

$$\eta = \sqrt{\frac{\mu}{\varepsilon}} = \sqrt{\frac{\mu_0}{\varepsilon_r \varepsilon_0}}$$
$$\sqrt{\varepsilon_r} = \frac{1}{\eta} \sqrt{\frac{\mu_0}{\varepsilon_0}} = 1.9986$$
$$\rightarrow \varepsilon_r = 4.0$$
(ii) Wave frequency:
$$\beta = 4\pi/3 = \omega \sqrt{\mu_0 \varepsilon}$$
$$\omega = \frac{4\pi}{3\sqrt{\mu_0 \varepsilon}}$$
$$2\pi f = 3.9946 \times 10^{16}$$

 $\rightarrow f = 99.93 \times 10^6 \approx (100 MHz)$

iii) Magnetic field equation:

$$\vec{H}(z,t) = \frac{377}{\eta} \cos(\omega t - \frac{4\pi}{3}z + \frac{\pi}{6})\hat{j} = 2\cos(\omega t - \frac{4\pi}{3}z + \frac{\pi}{6})\hat{j}$$

Example 9:

In a lossless medium for which $\eta = 60\pi$, $\mu_r = 1$ and $\mathbf{H} = -0.1 \cos (\omega t - z) \mathbf{a}_x + 0.5 \sin (\omega t - z) \mathbf{a}_y$ A/m, calculate ε_r , ω , and **E**.

Solution:

In this case , $\sigma=0$, $\alpha=0$, and $\beta=1$, so

$$\eta = \sqrt{\mu/\varepsilon} = \sqrt{\frac{\mu_o}{\varepsilon_o}} \sqrt{\frac{\mu_r}{\varepsilon_r}} = \frac{120\pi}{\sqrt{\varepsilon_r}} \quad \text{or} \quad \sqrt{\varepsilon_r} = \frac{120\pi}{\eta} = \frac{120\pi}{60\pi} = 2 \implies \varepsilon_r = 4$$
$$\beta = \omega \sqrt{\mu\varepsilon} = \omega \sqrt{\mu_o\varepsilon_o} \sqrt{\mu_r\varepsilon_r} = \frac{\omega}{c} \sqrt{4} = \frac{2\omega}{c}$$
$$\text{or} \qquad \omega = \frac{\beta c}{2} = \frac{1(3 \times 10^8)}{2} = 1.5 \times 10^8 \text{ rad/s}$$

$$\nabla \times \mathbf{H} = \sigma \mathbf{E} + \varepsilon \frac{\partial \mathbf{E}}{\partial t} \to \mathbf{E} = \frac{1}{\varepsilon} \int \nabla \times \mathbf{H} \, dt$$

where $\sigma = 0$.

But
$$\nabla \times \mathbf{H} = \begin{vmatrix} a_x & a_y & a_z \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial x} & \frac{\partial}{\partial z} \\ \mathbf{H}_x(z) & \mathbf{H}_x(z) & 0 \end{vmatrix} = -\frac{\partial \mathbf{H}_y}{\partial z} \mathbf{a}_x + \frac{\partial \mathbf{H}_x}{\partial z} \mathbf{a}_y$$

= 0.5 cos (
$$\omega$$
t - z) \mathbf{a}_x - 0.1 sin (ω t - z) \mathbf{a}_y

Hence

$$\mathbf{E} = \frac{1}{\varepsilon} \int \nabla \times \mathbf{H} dt = \frac{0.5}{\varepsilon \omega} \sin(\omega t - z) \mathbf{a}_x + \frac{0.1}{\varepsilon \omega} \sin(\omega t - z) \mathbf{a}_y$$

= 94.25 sin (\overline{\overline{t}} - z) \mathbf{a}_x + 18.85 cos (\overline{\overline{t}} - z) \mathbf{a}_y V/m

Example10:

A uniform plane wave propagating in a medium has $E = 2e^{-\alpha z} \sin(10^8 t - \beta z) \mathbf{a}_y$ V/m. If the medium is characterized by $\varepsilon_r = 1, \mu_r = 20$, and $\sigma = 3$ mhos/m, find α, β , and **H**.

Solution:

We need to determine the loss tangent to be able to tell whether the medium is a lossy dielectric or a good conductor.

$$\frac{\sigma}{\omega\varepsilon} = \frac{3}{10^8 \times 1 \times \frac{10^{-9}}{36\pi}} = 3393 >> 1$$

showing that the medium may be regarded as a good conductor at the frequency

of operation. Hence,
$$\alpha = \beta = \sqrt{\frac{\mu\omega\sigma}{2}} = \left[\frac{4\pi \times 10^{-7} \times 20(10^8)(3)}{2}\right]^{1/2} = 61.4$$

 $\alpha = 61.4 \text{ Np/m}$, $\beta = 61.4 \text{ rad/m}$
Also, $|\eta| = \sqrt{\frac{\mu\omega}{\sigma}} = \left[\frac{4\pi \times 10^{-7} \times 20(10^8)}{3}\right]^{1/2} = \sqrt{\frac{800\pi}{2}}$

$$\tan 2\theta_{\eta} = \frac{\sigma}{\omega\varepsilon} = 3393 \rightarrow \theta_{\eta} = 45^{\circ} = \frac{\pi}{4}$$

Hence

$$\mathbf{H} = H_o e^{-\alpha z} \sin\left(\omega t - \beta z - \frac{\pi}{4}\right) \mathbf{a}_H$$

$$\mathbf{a}_{H} = \mathbf{a}_{k} \times \mathbf{a}_{E} = \mathbf{a}_{z} \times \mathbf{a}_{y} = -\mathbf{a}_{x}$$

and

$$H_o = \frac{E_o}{|\eta|} = 2\sqrt{\frac{3}{800\pi}} = 69.1 \times 10^{-3}$$

Thus
$$\mathbf{H} = -69.1 \times 10^{-3} e^{-61.4z} \sin\left(10^8 t - 61.42 z - \frac{\pi}{4}\right) \mathbf{a}_x \text{ mA/m}$$

Example11:



Power of electricity generated = 63.2 mW

References:

- Justin Peatross, & Michael Ware, Physics of Light and Optics, Brigham Young University, Edition July 20, 2015, Justin Peatross and Michael Ware optics.byu.edu.2015
- 2. 1000 Solved Problems in Classical Physics. Ahmad A. Kamal. 2011.
- Justin Peatross, & Michael Ware, Physics of Light and Optics, Brigham Young University, Edition July 20, 2015, Justin Peatross and Michael Ware optics.byu.edu.2015